

Multi-pion Bose-Einstein correlations and the search for quantum coherence

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GDRE 2014

Microscopy (10^{-6})

Femtoscscopy (10^{-15})

Femtoscscopy = The study of particle correlations at low relative momentum

In high multiplicity events ($M > \sim 50$) such correlations are dominated by:
Bose-Einstein correlations (Quantum Statistics)
and
Final-State-Interactions (Coulomb+strong)

Only 2 sources of correlations make femtoscscopy a clean probe of heavy-ion physics.

Use Measuring the source radius

1

$$\sim \Delta x$$

- The last stage of particle interactions is kinetic freeze out.
- At freeze out in high-energy particle collisions, the characteristic separation of particles is femtoscopic ($\Delta x \sim 10^{-15}$ m).

$$\Delta x \Delta p \gg 2\pi\hbar$$

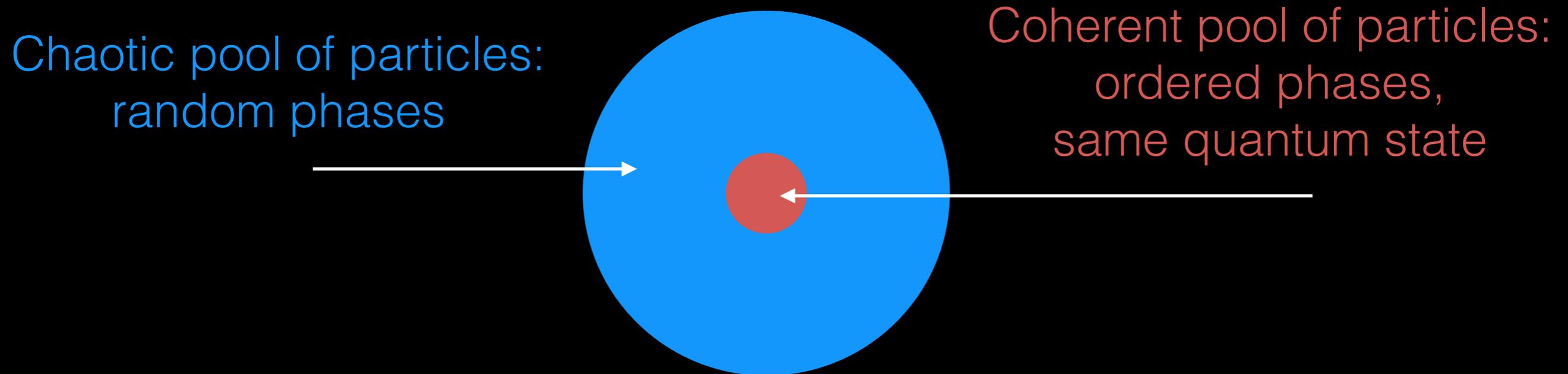
Classical:
no observable quantum phenomena

$$\Delta x \Delta p \sim 2\pi\hbar$$

Non Classical:
Bose-Einstein / Fermi-Dirac correlations

- Bose-Einstein correlations will be visible for $\Delta p < \sim 0.5$ GeV/c.
Relative momentum correlations are sensitive to the relative separation at freeze out.

Use 2 Measuring the fraction of pions which are coherent



Pion condensation, Disoriented Chiral Condensates, +.....
may create a coherent pool of pions.

Standard Correlation Functions

$$C_n = \frac{N_n(\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_n)}{N_1(\mathbf{p}_1)N_1(\mathbf{p}_2)\dots N_1(\mathbf{p}_n)}$$

\mathbf{p} = momentum

Projection Variables

$$q_{ij} = \sqrt{-(p_i - p_j)_\mu (p_i - p_j)^\mu}$$

$$k_T = |\vec{p}_{T1} + \vec{p}_{T2}|/2$$

$$Q_3 = \sqrt{q_{12}^2 + q_{13}^2 + q_{23}^2}$$

$$K_{T,3} = |\vec{p}_{T1} + \vec{p}_{T2} + \vec{p}_{T3}|/3$$

$$Q_4 = \sqrt{q_{12}^2 + q_{13}^2 + q_{14}^2 + q_{23}^2 + q_{24}^2 + q_{34}^2}$$

$$K_{T,4} = |\vec{p}_{T1} + \vec{p}_{T2} + \vec{p}_{T3} + \vec{p}_{T4}|/4$$

Fitting 2- and 3-pion Bose-Einstein Correlations

$$C_2^{\text{QS}}(q) = 1 + \lambda E_w^2(R_{\text{inv}} q) e^{-R_{\text{inv}}^2 q^2}$$

Quantum-Statistics
correlation

Suppression parameter
related to incorrect fit functions
and coherence

$$E_w(R_{\text{inv}} q) = 1 + \sum_{n=3}^{\infty} \frac{\kappa_n}{n!(\sqrt{2})^n} H_n(R_{\text{inv}} q)$$

Edgeworth expansion for
non-Gaussian features

Isolation and fitting of 3-pion Bose-Einstein Cumulant

$c_3 =$ 3-pion cumulant correlation

Isolated using:

$$N_3(p_1, p_2, p_3)$$

— 3 pions from same event

$$N_2(p_1, p_2)N_1(p_3)$$

— 2 pions from same event

$$N_1(p_1)N_1(p_2)N_1(p_3)$$

— All pions from different events

$$K_3 = K_2^{12} * K_2^{13} * K_2^{23}$$

— 3-body Final-State-Interaction

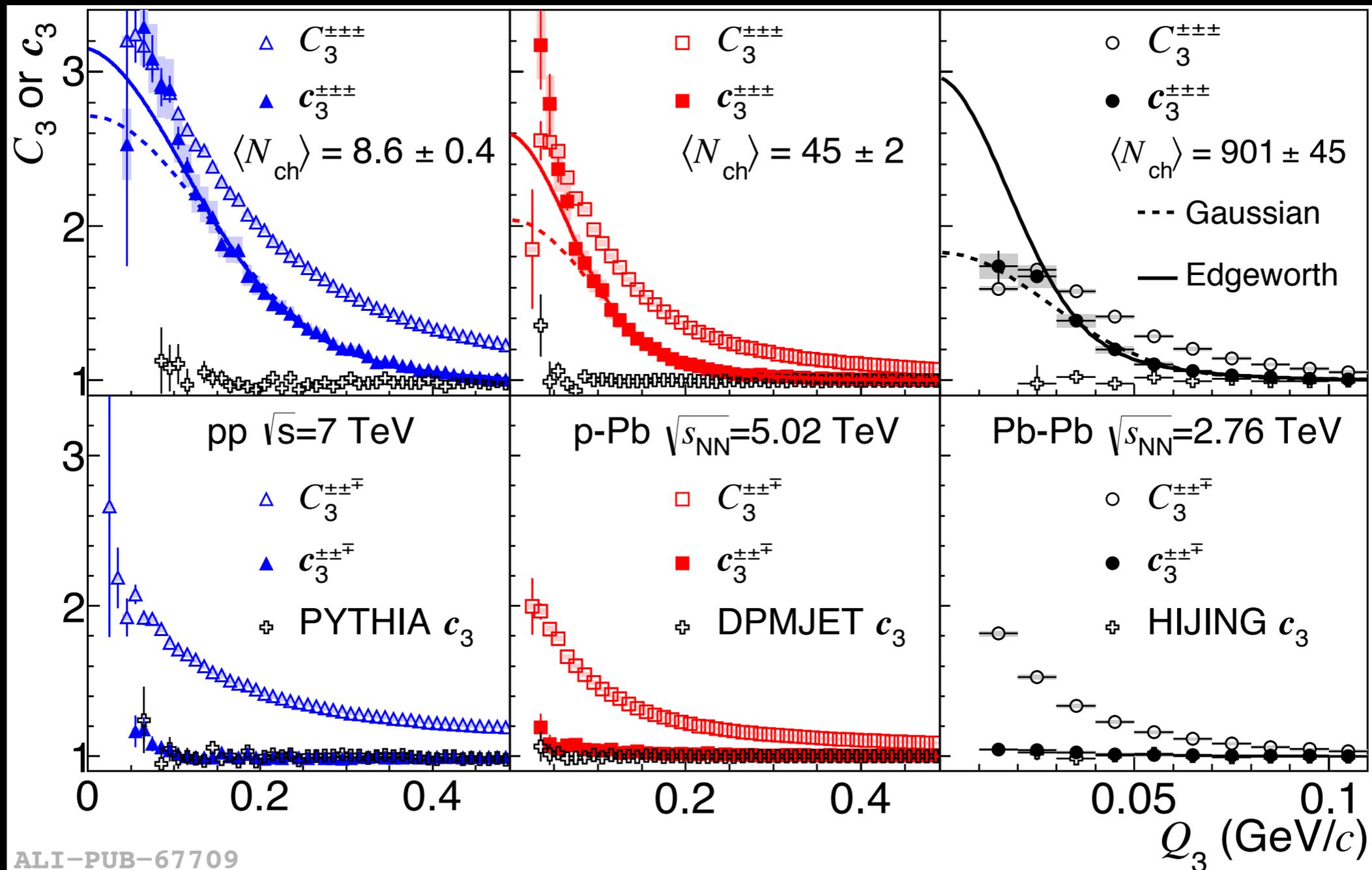
Cumulant fitting

$$c_3(q_{12}, q_{31}, q_{23}) = 1 + \lambda_3 E_w(R_{inv,3} q_{12}) E_w(R_{inv,3} q_{31}) E_w(R_{inv,3} q_{23}) e^{-R_{inv,3}^2 Q_3^2/2}$$

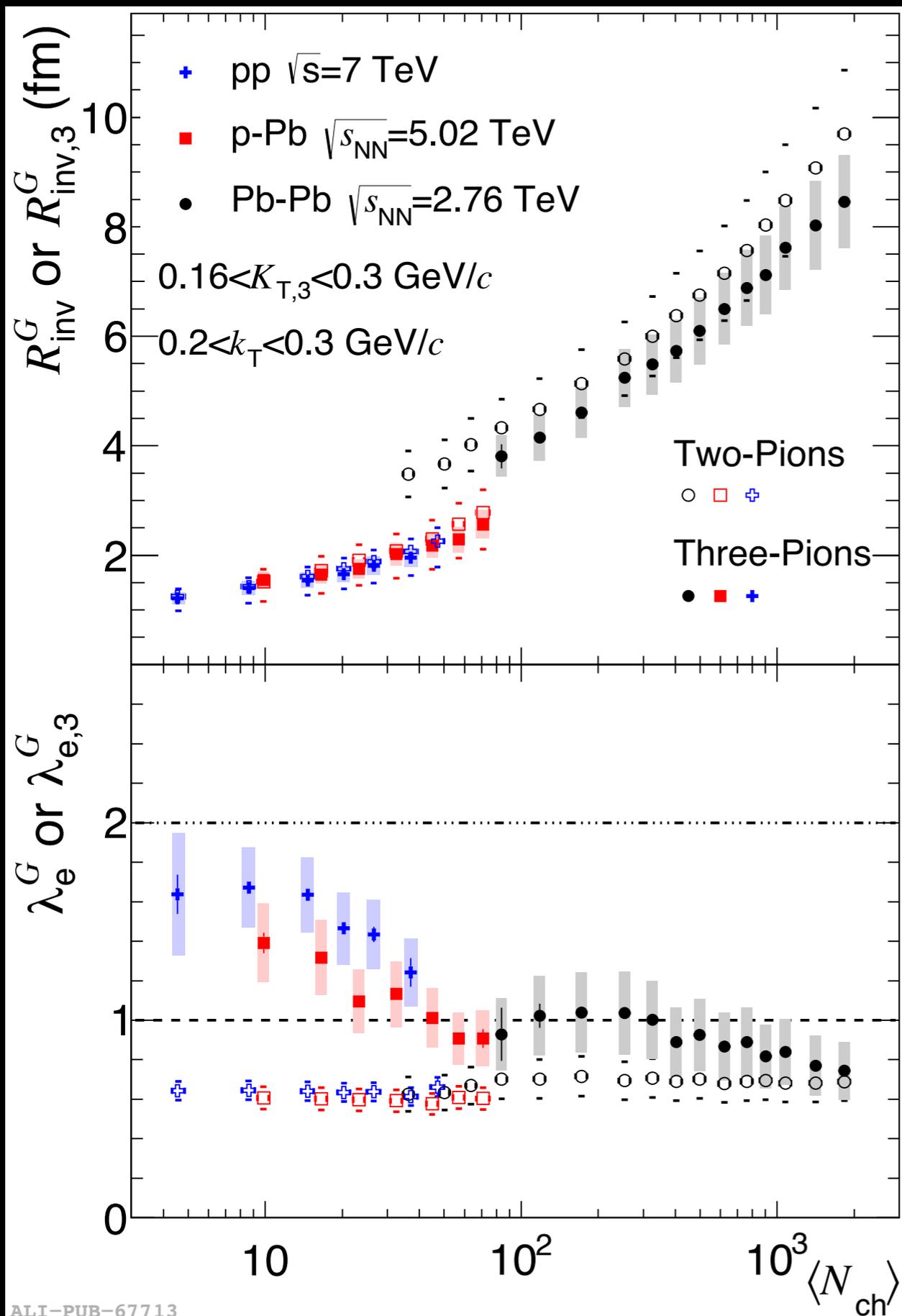
3-pion suppression parameter

Edgeworth factors to account for
non-Gaussian features

3-pion Femtoscopic Correlation Functions



ALI-PUB-67709

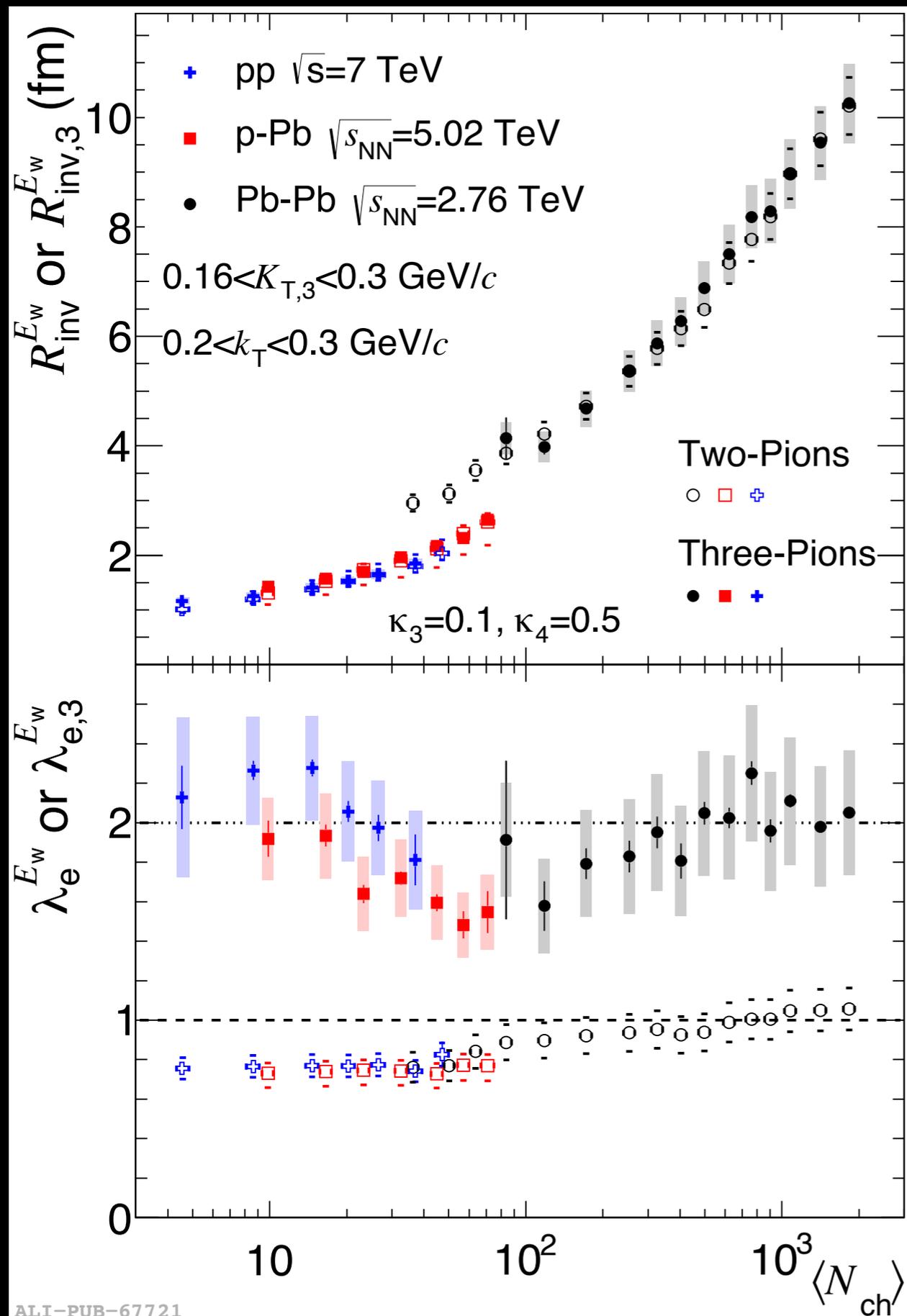


1D source radii from 2- and 3-pion correlations in pp, p-Pb and Pb-Pb

Gaussian fits

radii from 3-pion correlations somewhat smaller than from 2-pions.

Intercept parameters far below the chaotic limits.



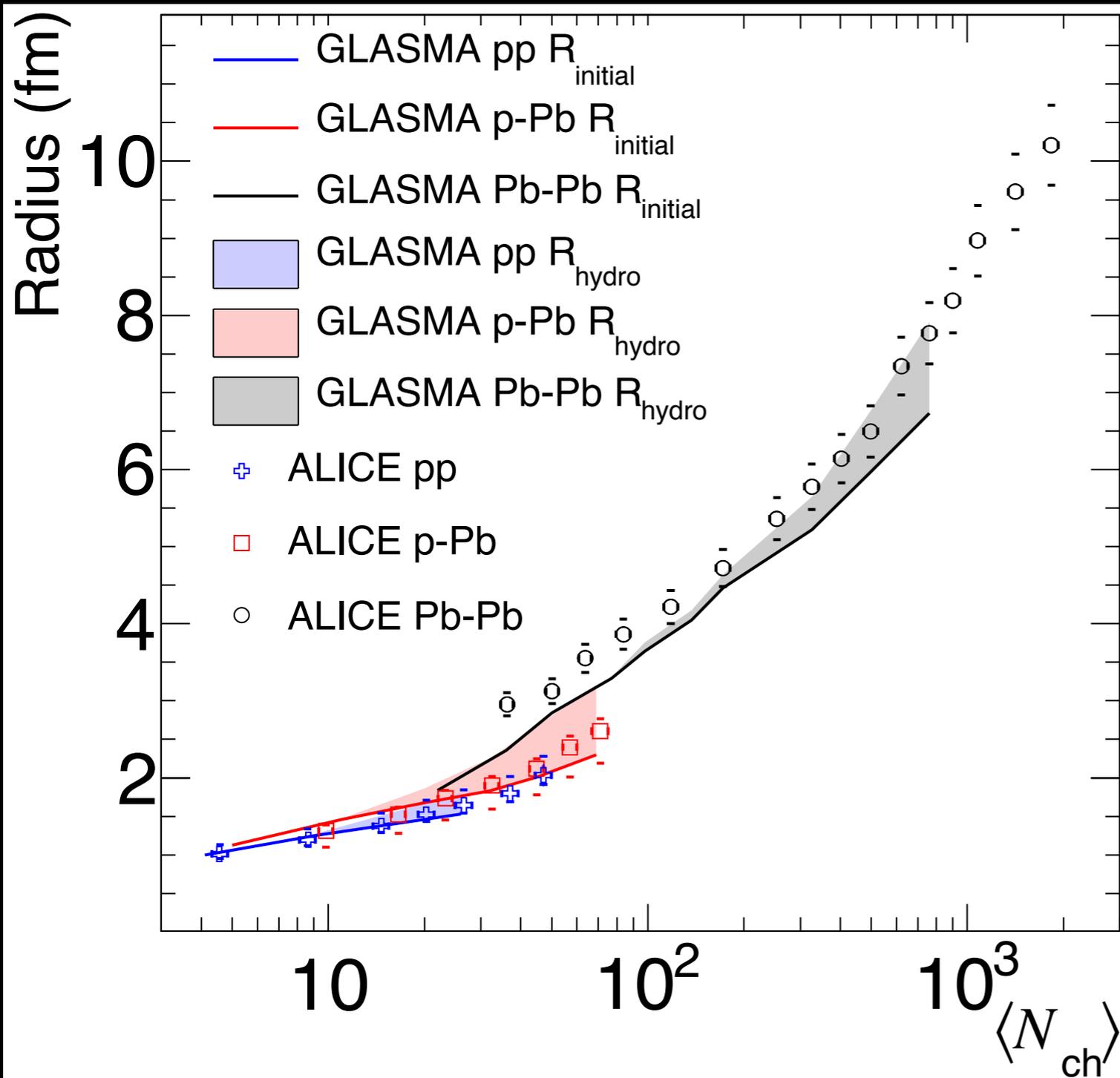
1D source radii from 2- and 3-pion correlations in pp, p-Pb and Pb-Pb

Edgeworth fits

radii from 3-pion correlations consistent with that from 2-pions.

Intercept parameters much more consistent with chaotic limits.

Radii Comparison with IP-GLASMA



Message:
Similarity of ALICE radii
in p-Pb and pp
can be reproduced with GLASMA
initial conditions alone.

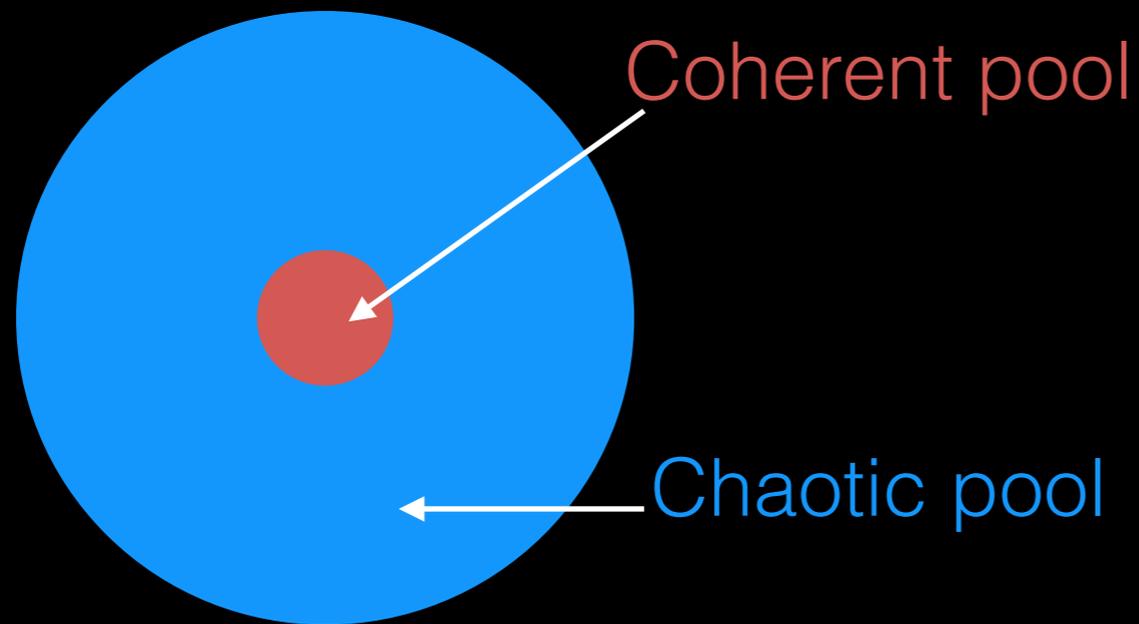
They can also be reproduced
with a hydrodynamic phase
in p-Pb.

Schenke & Venugopalan
arXiv:1405.3605

Use
2

Measuring the coherent
fraction of pions

2-pion Bose-Einstein contributions



2-pions

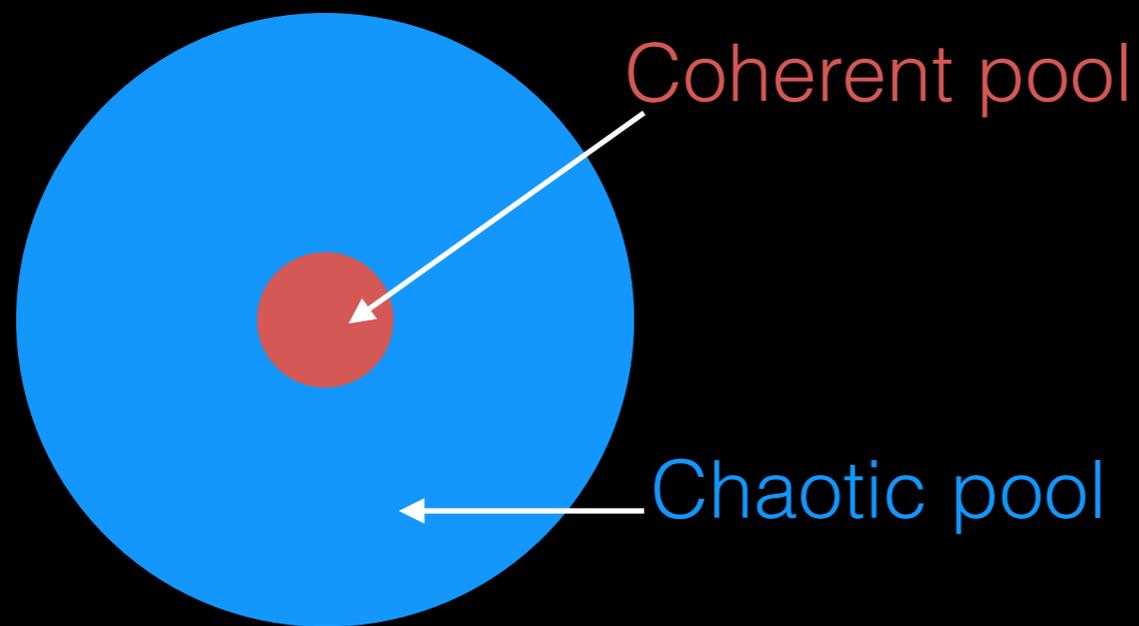
$\pi \pi$

$\pi \pi$

~~$\pi \pi$~~

1 suppressed combination

3-pion Bose-Einstein contributions



3-pions

$\pi \pi \pi$

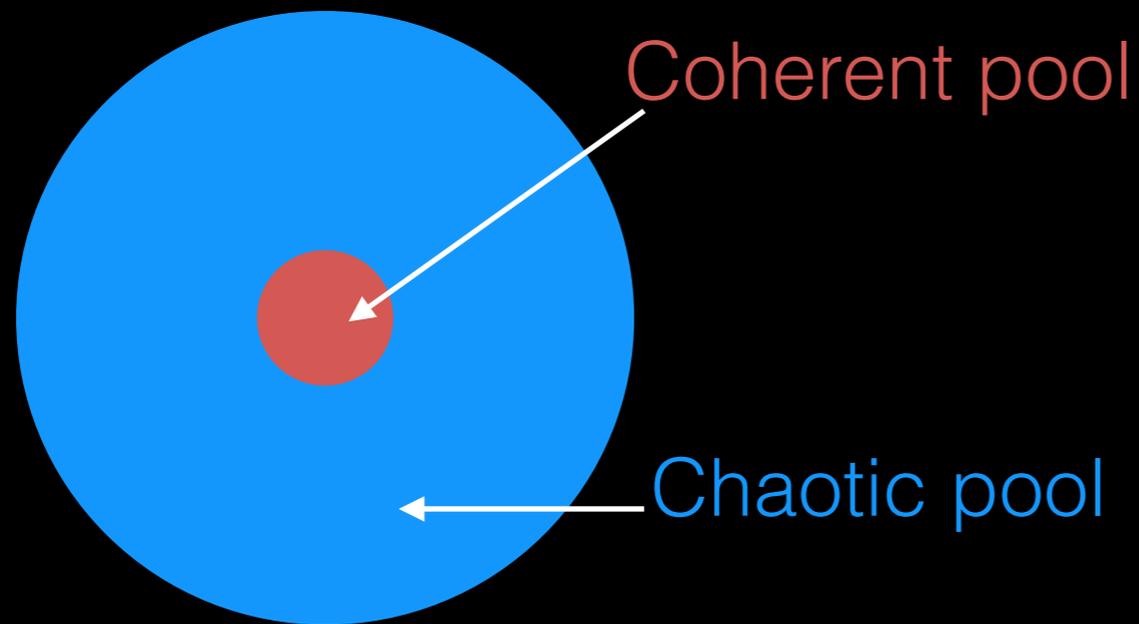
$\pi \pi \pi$

~~$\pi \pi \pi$~~

~~$\pi \pi \pi$~~

2 suppressed combinations

4-pion Bose-Einstein contributions



4-pions

$\pi \pi \pi \pi$

3 suppressed combinations

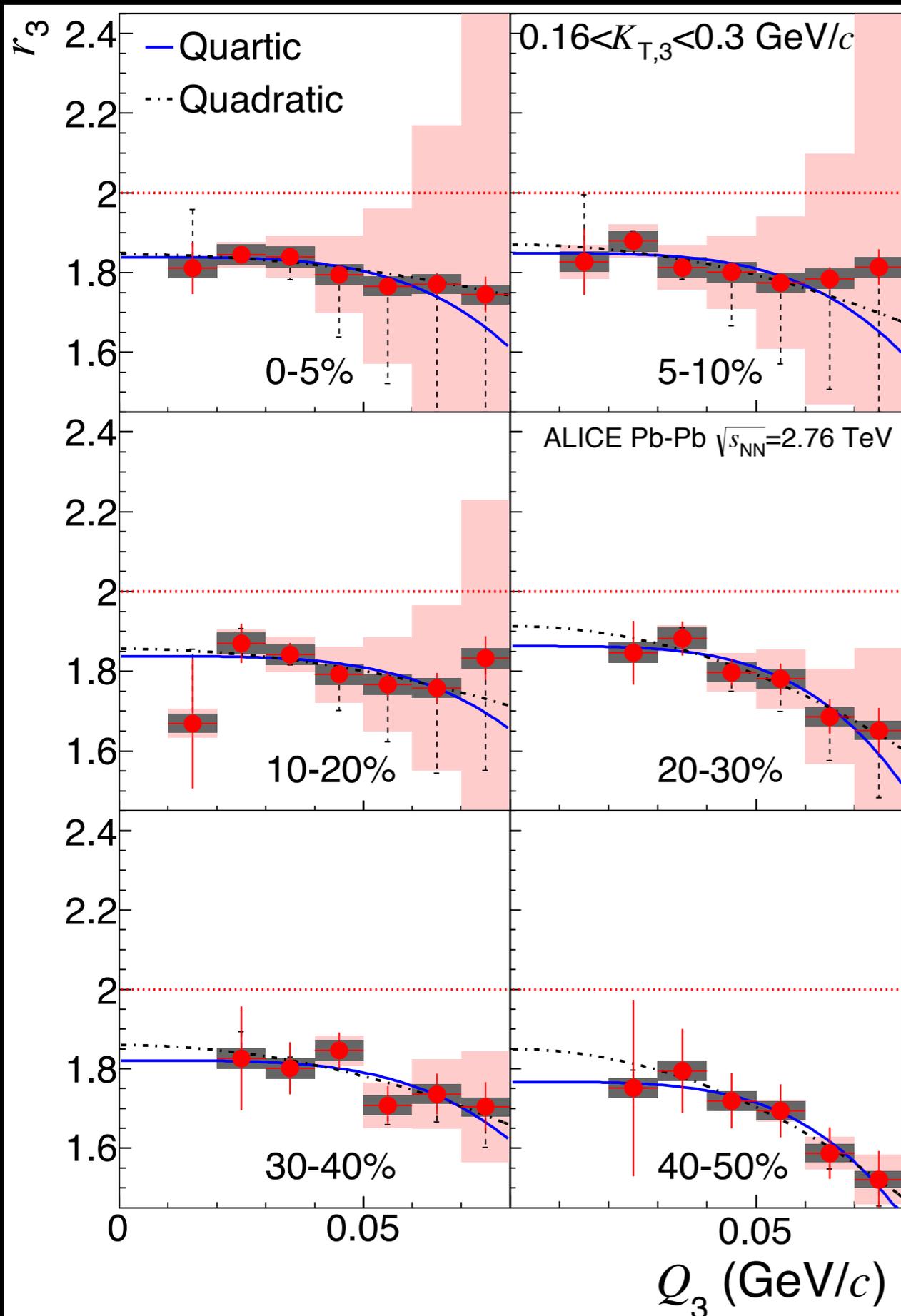
Resolution of coherence increases with the number of pions used.

Measure of coherent fraction
by comparing 3-pion to
2-pion correlation strength

$$r_3 = \frac{c_3(q_{12}, q_{23}, q_{31}) - 1}{\sqrt{(C_2(q_{12}) - 1)(C_2(q_{13}) - 1)(C_2(q_{23}) - 1)}}$$

All Correlations are first
Coulomb corrected.

r_3 is suppressed below 2.0.
Intercept corresponds to
~20% coherence at low p_T .



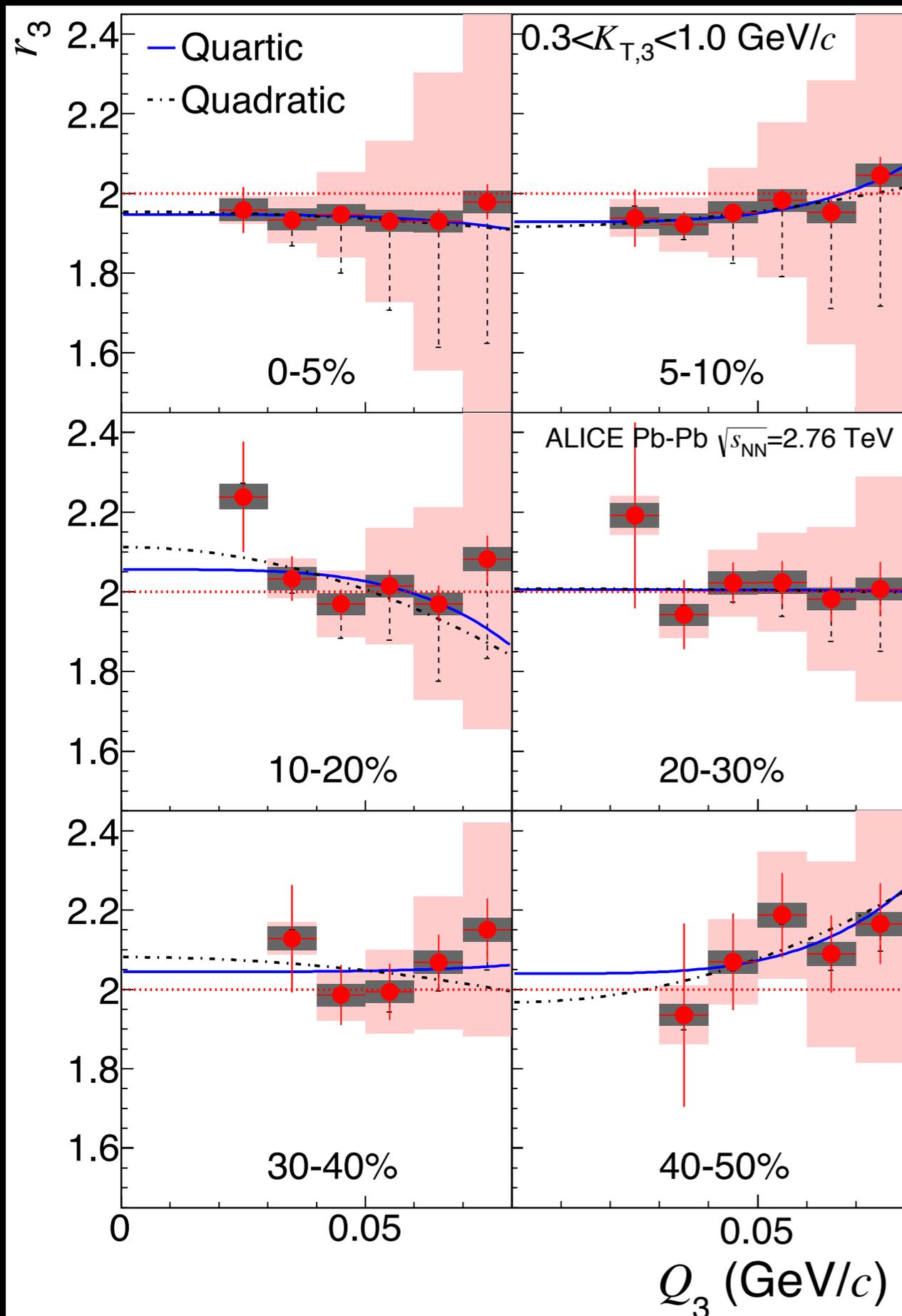
ALICE 2014
PRC 89 024911 (2014)

Measure of coherent fraction by comparing 3-pion to 2-pion correlation strength

$$r_3 = \frac{c_3(q_{12}, q_{23}, q_{31}) - 1}{\sqrt{(C_2(q_{12}) - 1)(C_2(q_{13}) - 1)(C_2(q_{23}) - 1)}}$$

All Correlations are first Coulomb corrected.

r_3 is consistent with 2.0.
Intercept is consistent with **0% coherence at high p_T**



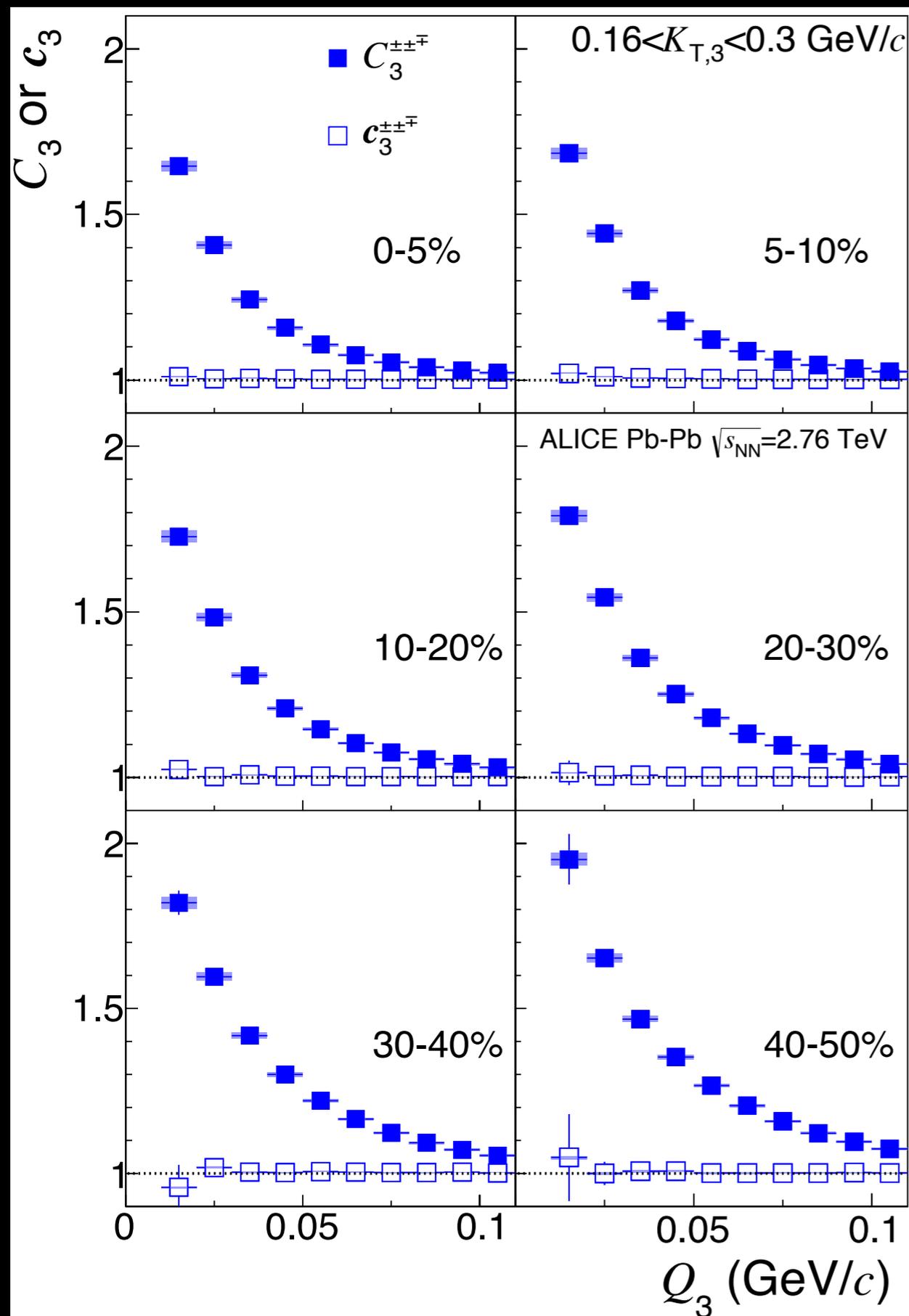
ALICE 2014
PRC 89 024911 (2014)

Check that 3-body Coulomb corrections work

The cumulant (hollow points) is Coulomb corrected.

Consistency with unity demonstrates success of 3-body Coulomb ansatz

$$K_3 = K_2^{12} * K_2^{13} * K_2^{23}$$



ALICE 2014
PRC 89 024911 (2014)

Future measurements in ALICE

An interesting hint of coherence in Pb-Pb collisions has been observed with 3-pion correlations.

4-pion correlations are even more sensitive to coherence. The 4-pion analysis is already underway. Preliminary results expected for WPCF Aug 2014.

New method to extract coherent fraction— “Built” C_4^{QS}

$C_2(k_T, q_{Out}, q_{Side}, q_{Long})$ is statistically well defined.
It can be tabulated very differentially.

For fully chaotic pion emission:

$$C_2^{QS}(q_{ij}) - 1.0 = T_{ij}^2$$

(pair exchange amplitude squared).

Also for fully chaotic emission and neglecting
the multi-pion phases for 3-pion and 4-pion interchanges,
 C_4^{QS} is entirely built from each of the 6 T_{ij}

Equations to Build QS correlations with coherence

G = coherent fraction of pions

$$C_2^{QS} - 1 = (1 - G^2)T_{12}^2 \quad (52)$$

$$C_3^{QS} - 1 = (1 - G)^2(T_{12}^2 + T_{13}^2 + T_{23}^2) \quad (53)$$

$$+ (6G(1 - G)^2 + 2(1 - G)^3)T_{12}T_{13}T_{23} \quad (54)$$

$$C_4^{QS} - 1 = (1 - G^2)(T_{12}^2 + T_{13}^2 + T_{14}^2 + T_{23}^2 + T_{24}^2 + T_{34}^2) \quad (55)$$

$$+ (4G(1 - G)^3 + (1 - G)^4(T_{12}^2T_{34}^2 + T_{13}^2T_{24}^2 + T_{14}^2T_{23}^2)) \quad (56)$$

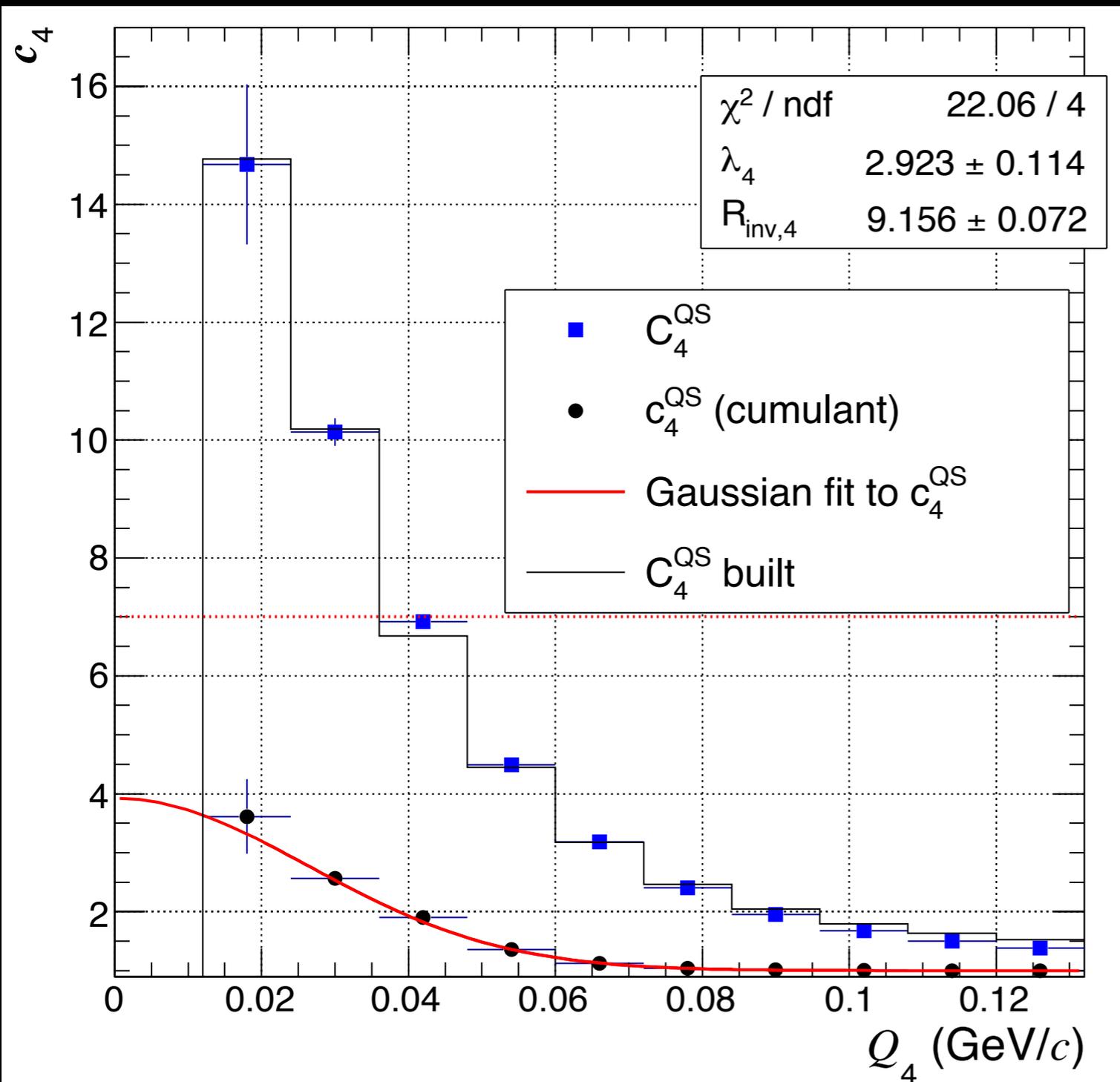
$$+ (6G(1 - G)^2 + 2(1 - G)^3)(T_{12}T_{13}T_{23} + T_{12}T_{14}T_{24} + T_{13}T_{14}T_{34} + T_{23}T_{24}T_{34}) \quad (57)$$

$$+ (8G(1 - G)^3 + 2(1 - G)^4)(T_{12}T_{13}T_{24}T_{34} + T_{12}T_{14}T_{23}T_{34} + T_{13}T_{14}T_{23}T_{24}) \quad (58)$$

These equations valid for $R_{coh}=R_{ch}$
(coherent Radius = chaotic Radius).

We will also consider $R_{coh}=0$ (point source).

Proof of Principle



Terminator model calculation (no coherence).

Black line (Built) is in good agreement with Blue points
→ 4-pion QS correlations are built from information already present at 2-pion level for the case of no coherence.

Deviations from this expectation may be used to search for non-zero coherence.

Summary

Radii Measurements ([Use 1](#))

For the first time, femtoscopic radii have been extracted using 3-pion cumulant Bose-Einstein correlations. 3-pion cumulants are advantageous in low-multiplicity events such as in pp and p-Pb collisions. They are less sensitive to background correlations.

Coherence Measurements ([Use 2](#))

An unexpected hint of quantum coherence has been observed using 3-pion Bose-Einstein correlations in Pb-Pb collisions at the LHC. The eventual 4-pion measurement should shed more light on the possibility of coherent pion emission.

Supporting Slides

Fitting 2-pion Bose-Einstein Correlations

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Quantum-Statistics
correlation

Suppression parameter
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$$E_w(R_{\text{inv}} q) = 1 + \sum_{n=3}^{\infty} \frac{\kappa_n}{n!(\sqrt{2})^n} H_n(R_{\text{inv}} q)$$

Edgeworth expansion for
non-Gaussian features

Gaussian Fit when $E_w = 1.0$

Non-Gaussian features parameterized with an Edgeworth expansion.

- κ_3 and κ_4 expansion parameters retained and extracted from 3-pion cumulants.
- Free fit performed for all 3 systems and all multiplicity bins.
- $\langle \kappa_3 \rangle = 0.1$
- $\langle \kappa_4 \rangle = 0.5$

Csorgo & Hegyi,
Phys. Lett. B 489, 15 (2000)

Isolation and fitting of 3-pion Bose-Einstein Cumulant

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Isolated using:

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— 3 pions from same event

$$N_2(p_1, p_2)N_1(p_3)$$

— 2 pions from same event

$$N_1(p_1)N_1(p_2)N_1(p_3)$$

— All pions from different events

$$K_3 = K_2^{12} * K_2^{13} * K_2^{23}$$

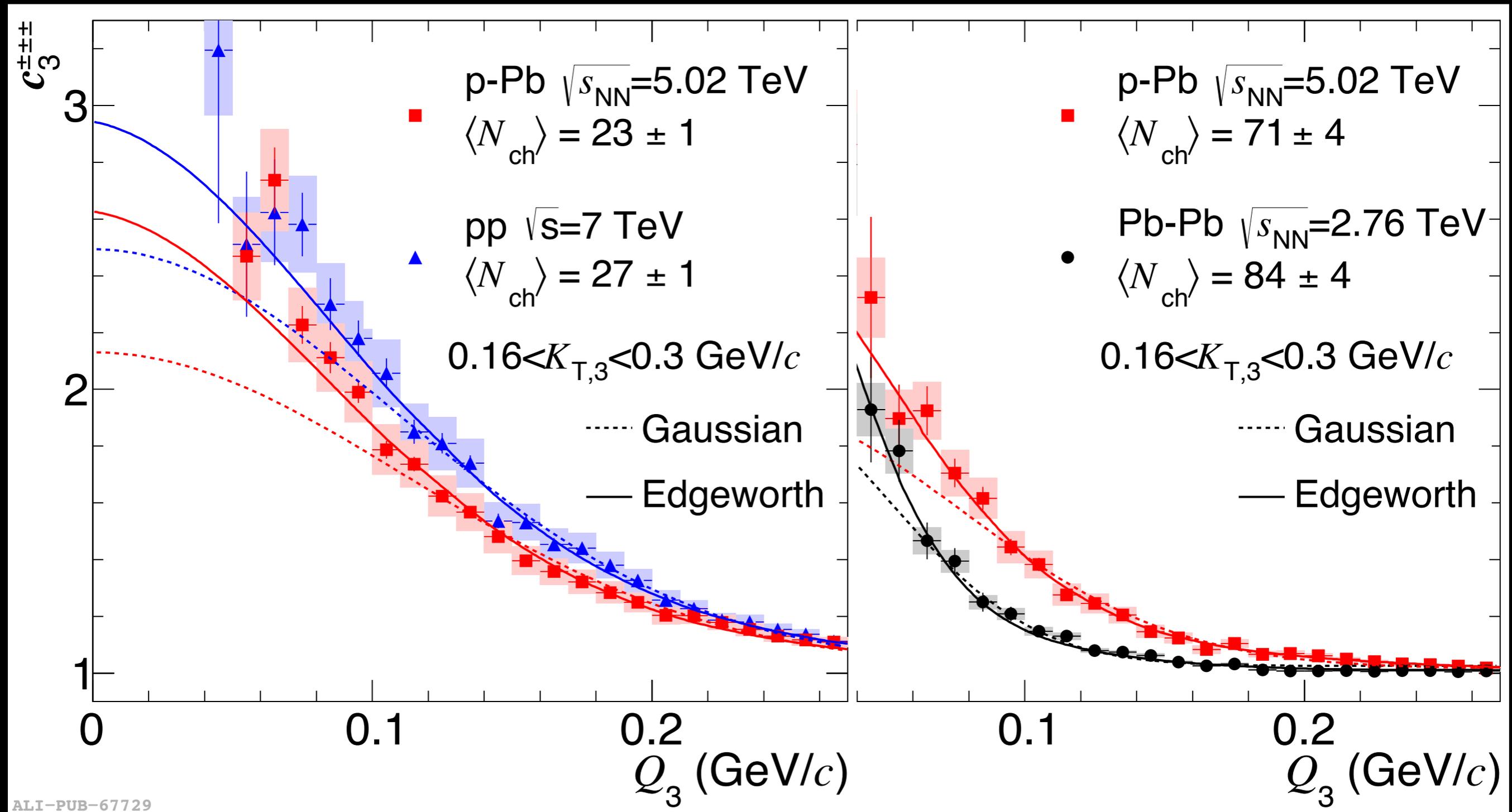
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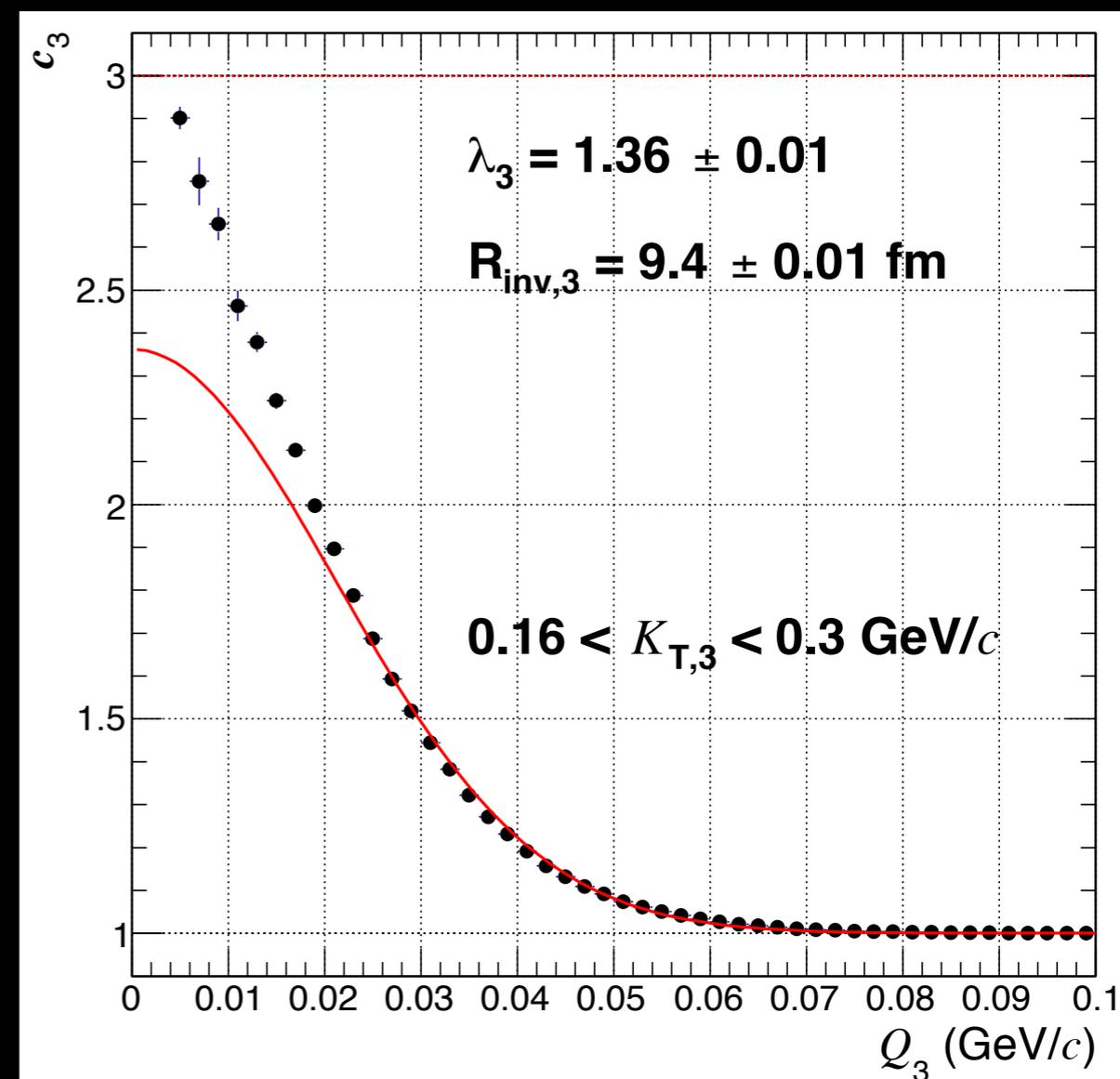
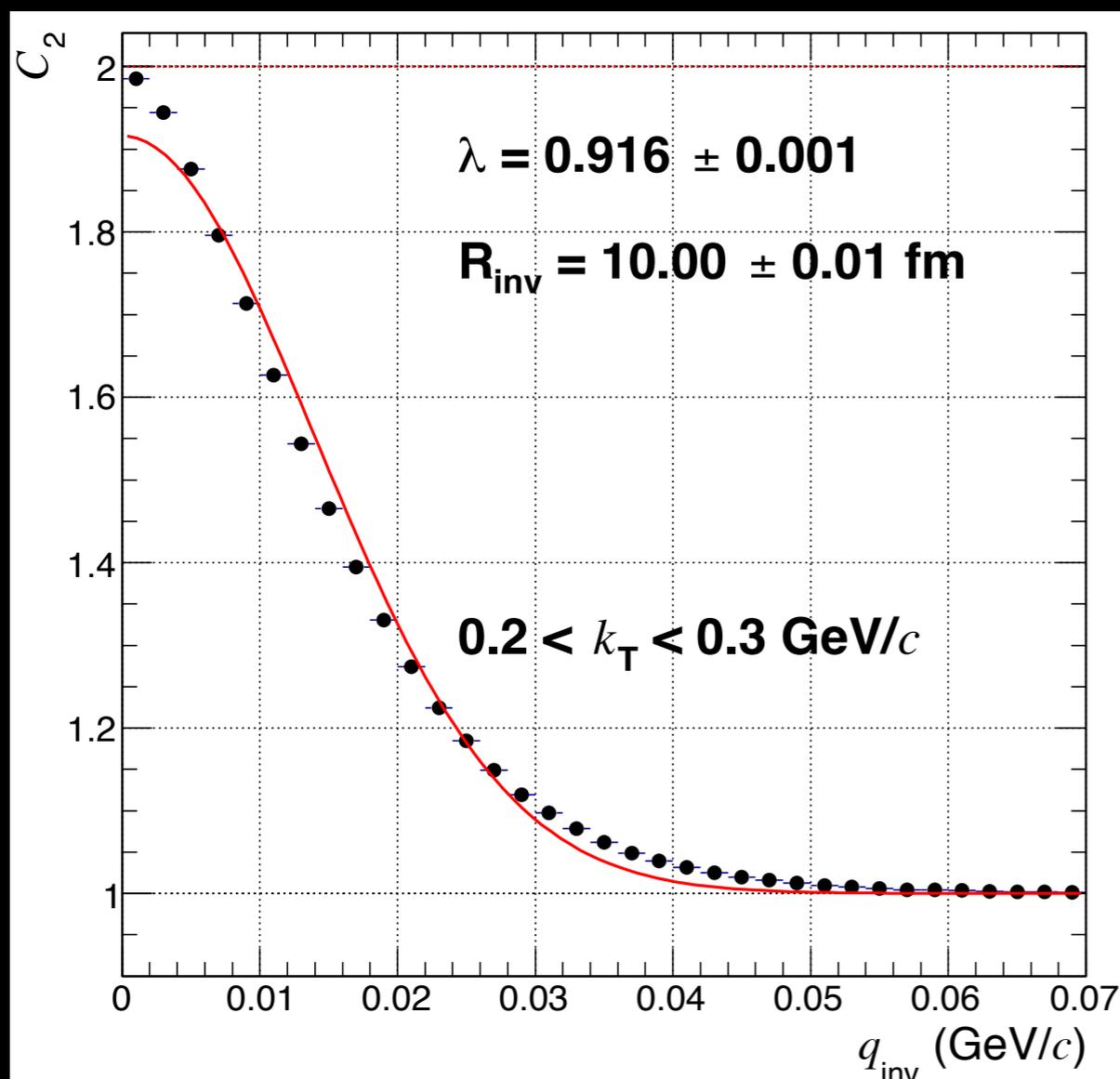
3-pion suppression parameter

Edgeworth factors to account for
non-Gaussian features



ALI-PUB-67729

Therminator2 calculations of 2- and 3-pion Bose-Einstein correlations

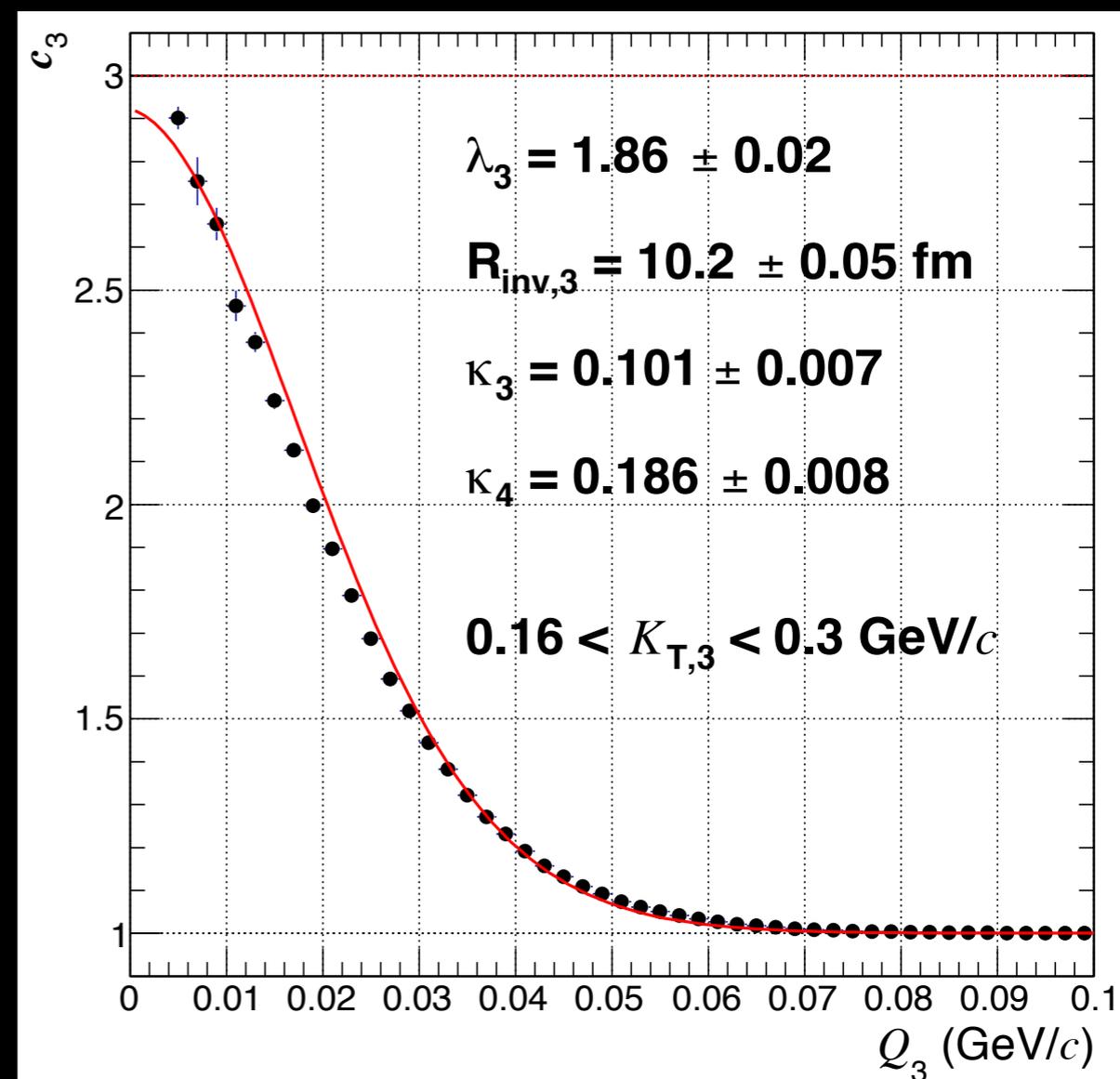
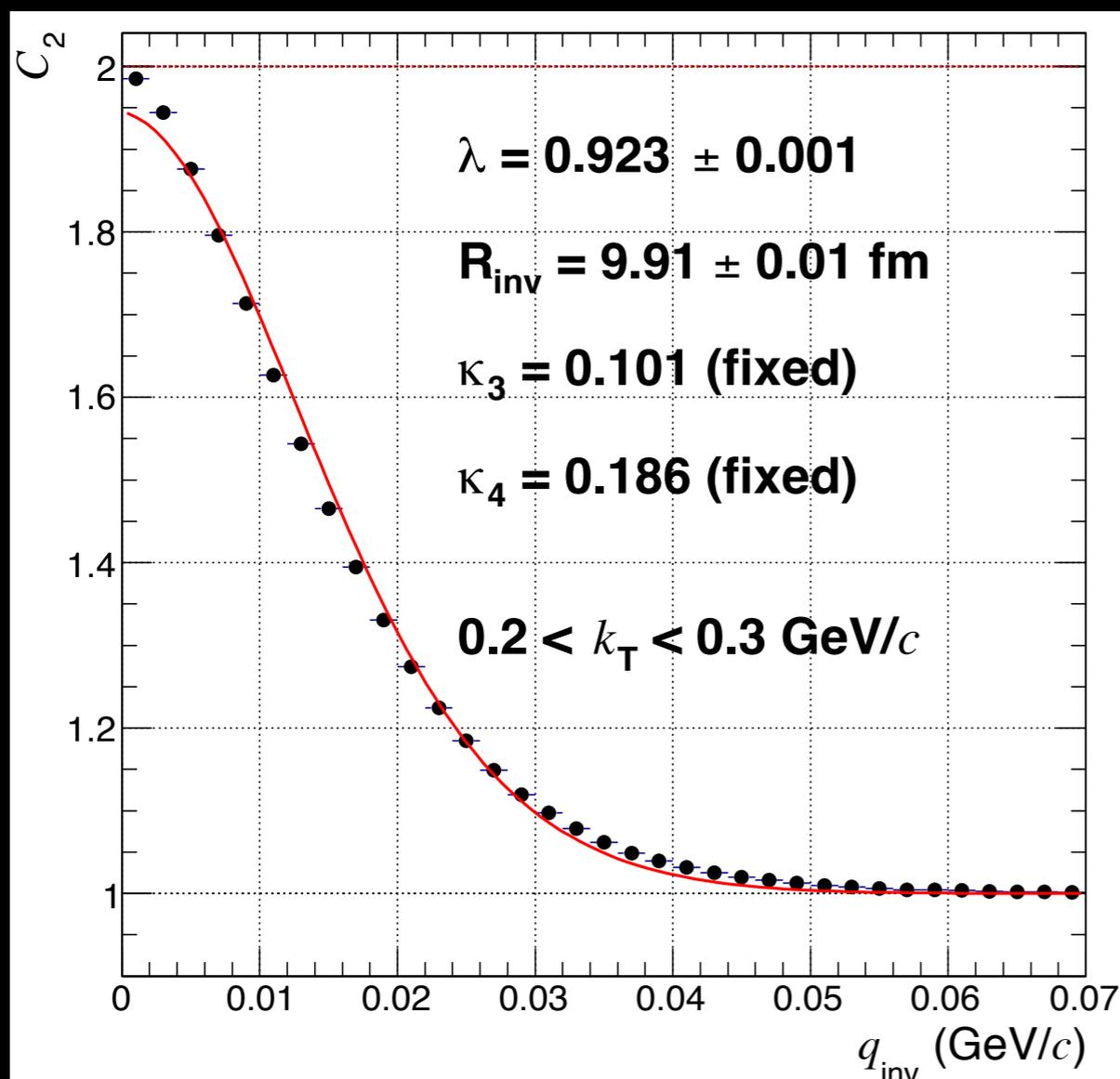


Gaussian fits in **red**.

$R_{inv,3}$ smaller than R_{inv} by $\sim 0.6 \text{ fm}$ (6%).

Therminator 2 (0-5% Pb-Pb)
Kisiel et al.,
Comput. Phys. Commun. 174, 669
(2006)
P. Bozek Phys. Rev. C 77, 034911
(2008)

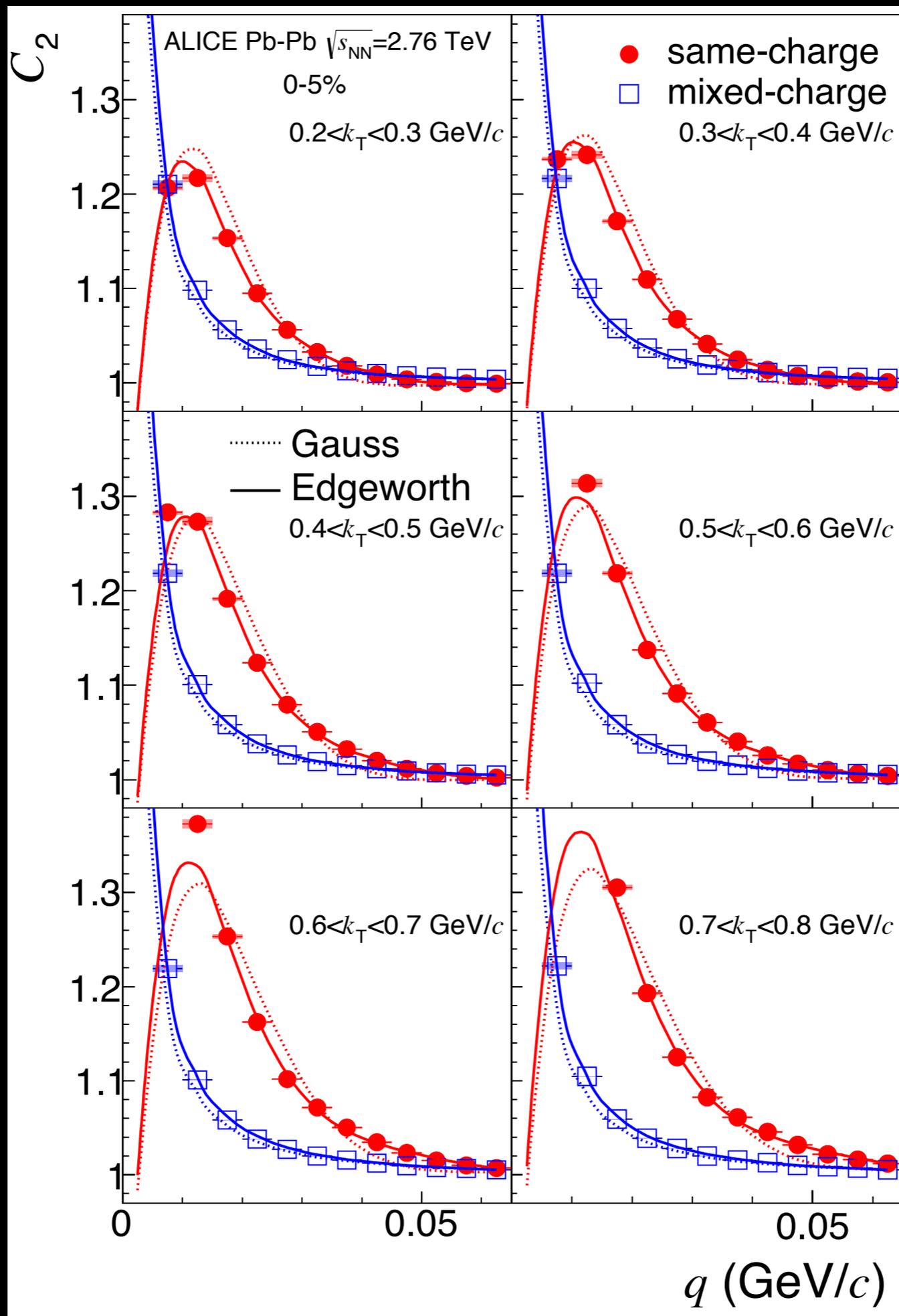
Therminator2 calculations of 2- and 3-pion Bose-Einstein correlations

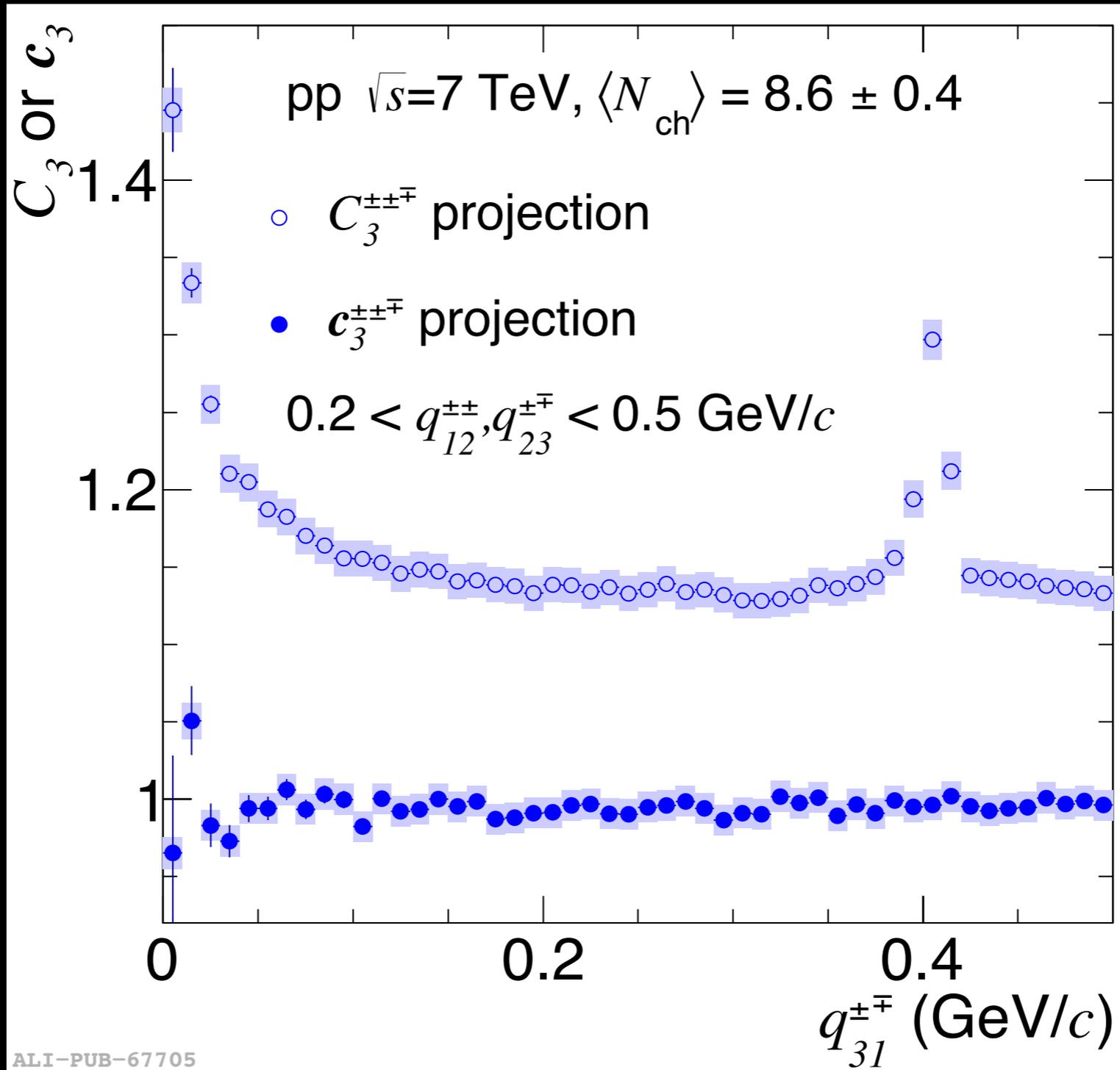


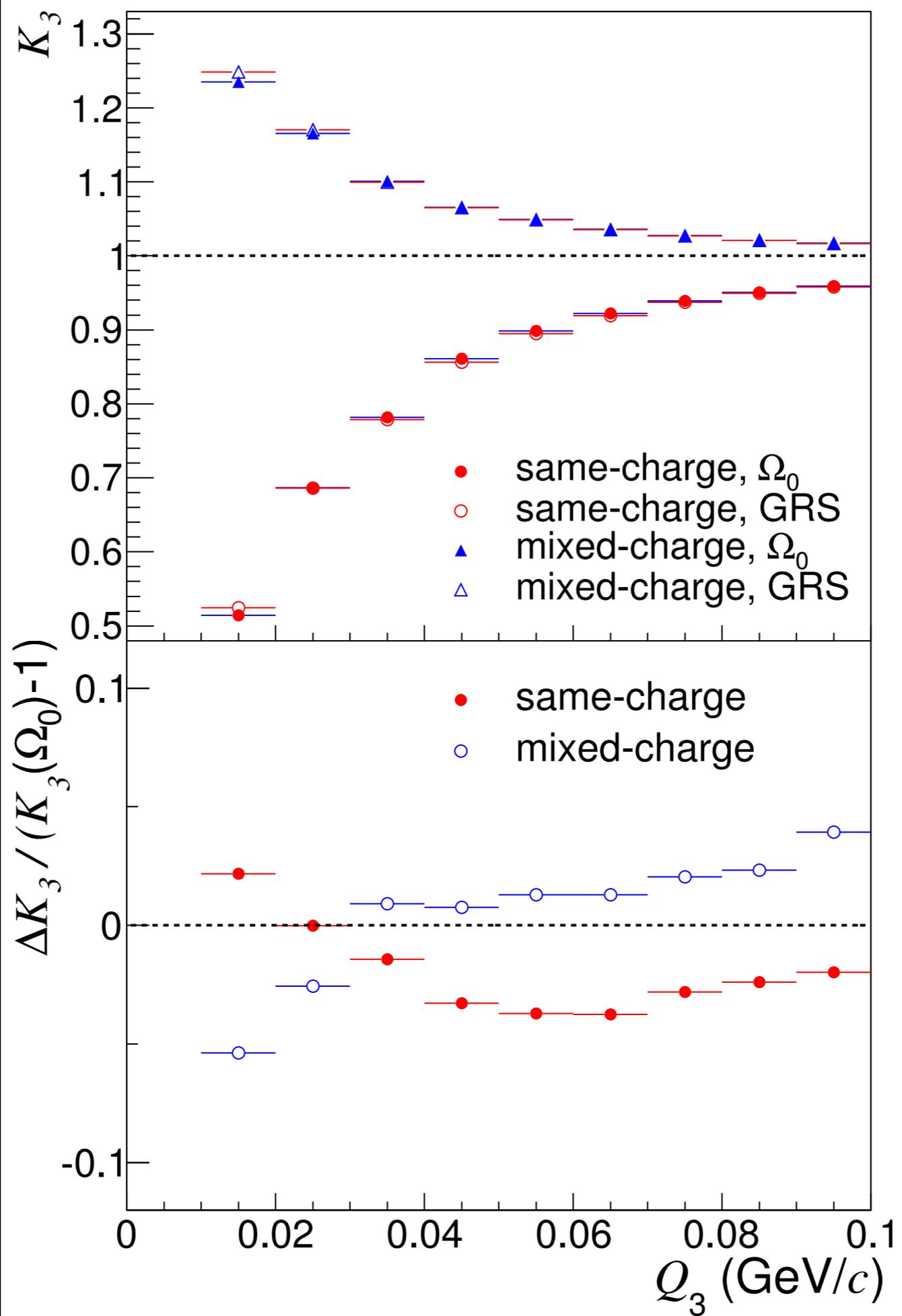
Edgeworth fits in **red**.

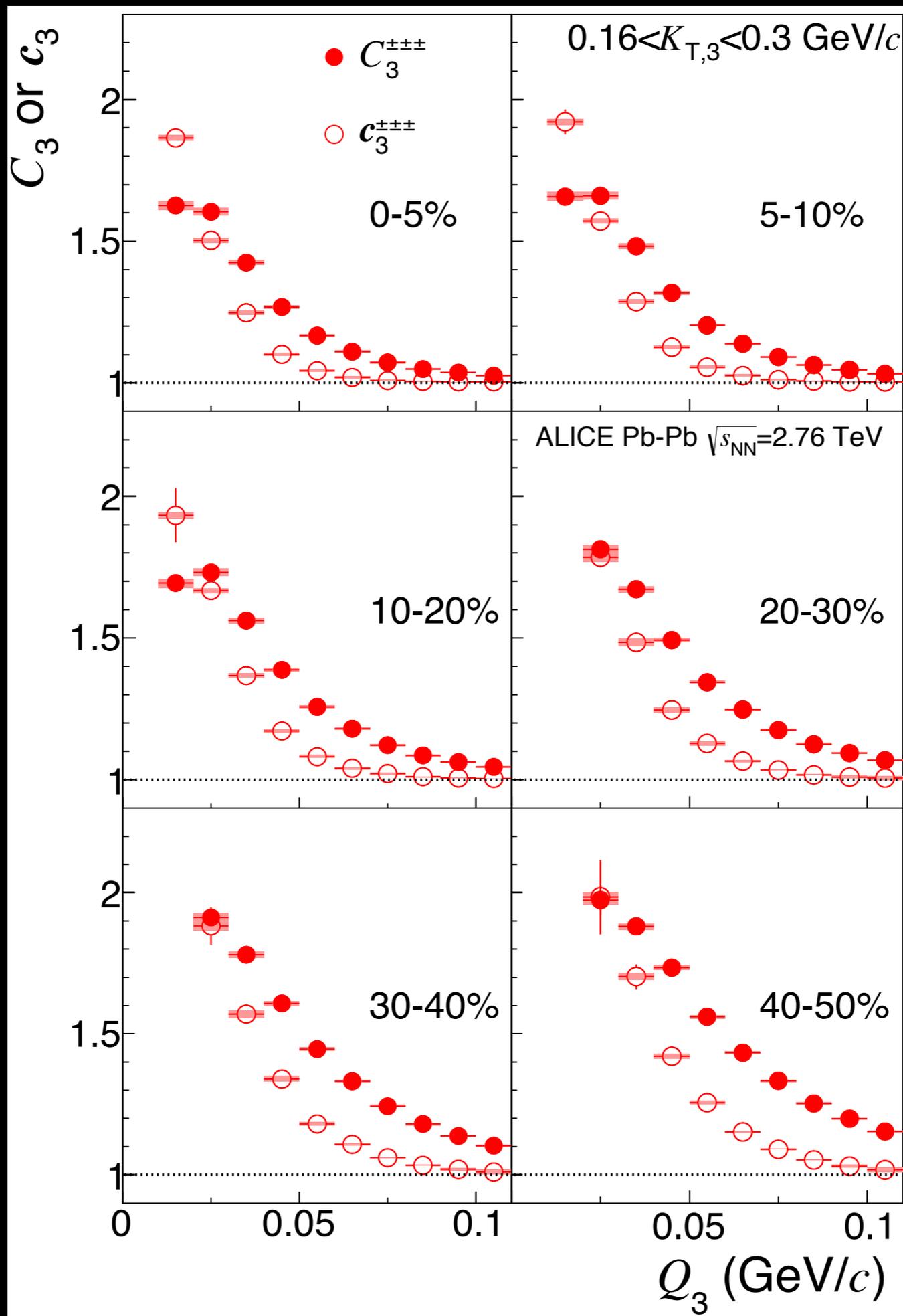
$R_{inv,3}$ similar to R_{inv} within $\sim 0.3 \text{ fm}$ (3%).

Therminator 2 (0-5% Pb-Pb)
Kisiel et al.,
Comput. Phys. Commun. 174, 669
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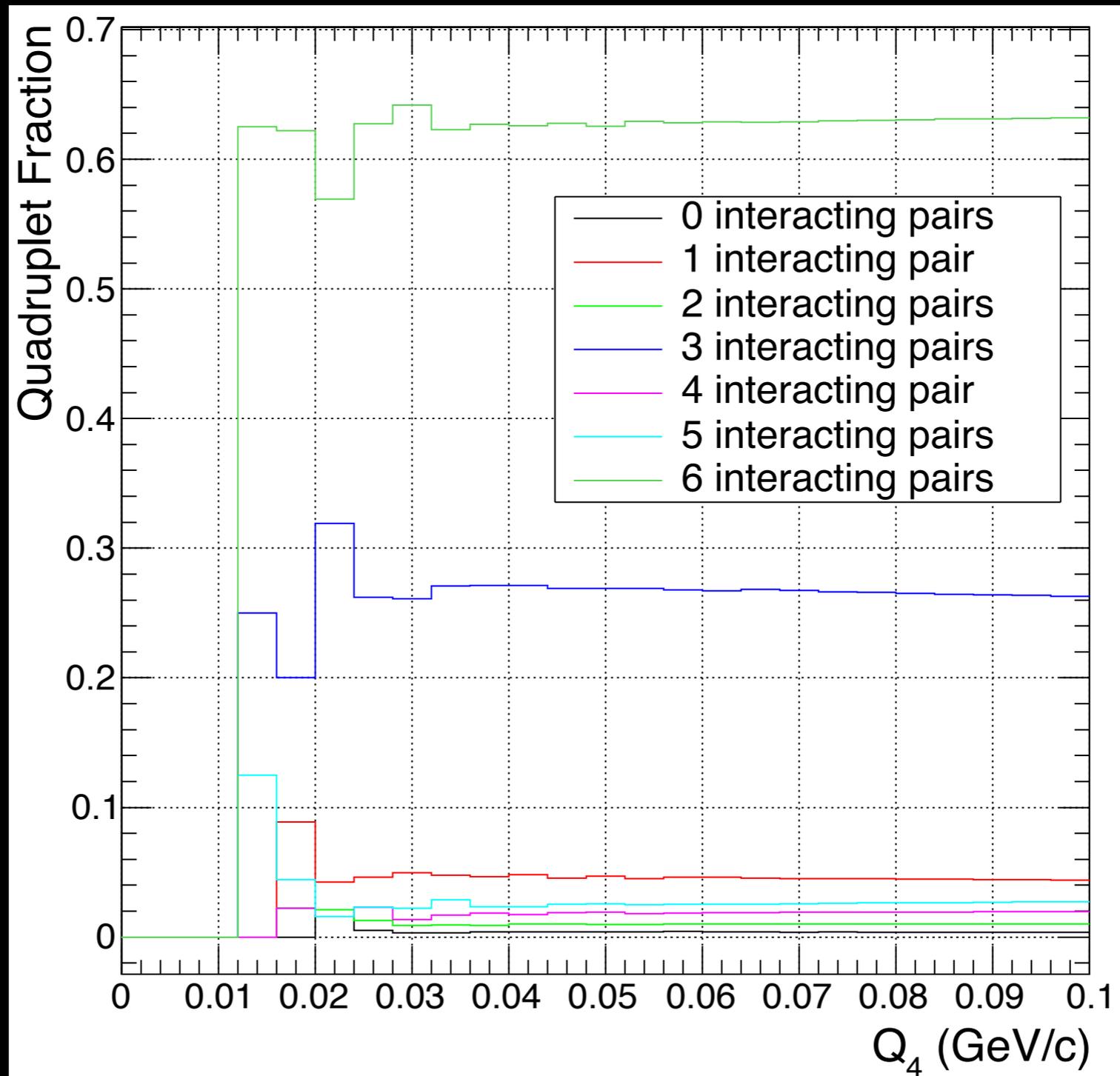








Quadruplet Fractions in Therminator



r_3 in Therminator

